

International Journal of Modern Physics A  
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## KAONIC HYDROGEN MEASUREMENT WITH DEAR AT DAΦNE

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Revised (Day Month Year)

The DEAR (DAΦNE Exotic Atom Research) experiment<sup>1</sup> measured the energy of X-rays emitted in the transitions to the ground states of kaonic hydrogen. The shift  $\epsilon$  and the width  $\Gamma$  of the  $1s$  state are sensitive quantities for tests of the current understanding of the low energy antikaon-nucleon interaction. We obtain  $\epsilon_{1s} = -193 \pm 37$  (stat.)  $\pm 6$  (syst.) and  $\Gamma_{1s} = 249 \pm 112$  (stat.)  $\pm 30$  (syst.)

*Keywords:* kaonic atoms; low-energy QCD.

## 1. Introduction

In the experiment X-rays emitted in the transitions to the ground states of kaonic atoms were measured.

A kaonic atom is formed when a negative kaon enters a medium, loses its kinetic energy through ionization and excitation of the atoms and molecules and eventually is captured, replacing the electron, in an excited orbit ( $n \simeq 25$ ). Via different cascade processes (Auger effect, Coulomb deexcitation, scattering) the kaonic atom deexcites to lower states.

When a kaon reaches a low- $n$  state with small angular momentum, strong interaction with the nucleus causes its absorption. The strong interaction is the reason for a shift in the energies of the low-lying levels from the purely electromagnetic values and the finite lifetime of the state corresponds to an increase in the observed level width.

Measuring shift and width allows to study the  $K^-p$  interaction at threshold using some counterpart of Deser-type relations<sup>2</sup>.

In a previous experiment<sup>3</sup> at KEK Iwasaki et al. have found that the strong interaction shift of the  $1s$  state in kaonic hydrogen is repulsive in agreement with scattering data and in disagreement with earlier kaonic hydrogen experiments. The origin of this repulsive strong interaction at threshold can be traced back to the presence of the  $\Lambda(1405)$  resonance which leads on one hand to a repulsive  $K^-p$  scattering length and on the other to the possible existence of strongly bound kaonic states in light nuclei<sup>4</sup>.

For a quantitative test of recent theoretical work<sup>5</sup> on the topic, additional data with better precision and finally a kaonic deuterium measurement is required. With DEAR we succeeded to perform an improved measurement of kaonic hydrogen. An experiment on kaonic deuterium is currently under development – the SIDDHARTA project.

Finally we want to point out that the scattering lengths are related to the so called meson-nucleon sigma terms<sup>6</sup> as shown by 'low energy theorems'. These terms are important quantities in non perturbative Quantum Chromo Dynamics (QCD) since they provide a measure of the chiral symmetry breaking.

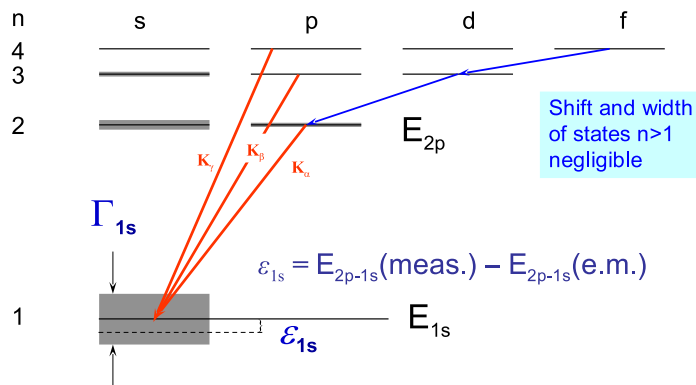


Fig. 1. Schematic view of the energy levels of kaonic hydrogen showing shift and width

## 2. Experiment

Our experiment was performed at the electron-positron collider DAΦNE at Frascati and used for the first time low energy, nearly monoenergetic kaons from the decay of  $\Phi$ s produced almost at rest for the formation of kaonic hydrogen. The small energy spread of the kaons introduced by a lateral boost of 15.2 MeV/c due the crossing angle of  $e^+ e^-$  was compensated by a special degrader. The use of a lepton machine results in a very low pion background as compared to proton machines. The low energy kaons were stopped in a special cryogenic light weight hydrogen gas target consisting of specially chosen structure materials to suppress fluorescence X ray radiation as far as possible. CCDs (Marconi Applied Technologies, CCD55-30) were used as X-ray detectors. Their advantages are large area, pixel analysis capability for background reduction and excellent energy resolution.

The target was operated with hydrogen at 25 K and 2 bar resulting in a density of 2.1 g/l (0.03 of liquid hydrogen density). 16 CCD detectors with a total area of 116cm<sup>2</sup> and a resolution of 160eV FWHM at 7.6 keV were used. The 10<sup>6</sup> pixels of each chip were read out every 2 min. The intensity of the charged kaons was measured by two plastic scintillators which detected the back to back emitted kaons<sup>7</sup>. The kaons initial energy of  $\approx 16$  MeV was reduced by plastic degraders to  $\approx 3$  MeV which allows to stop them within a few cm in the hydrogen gas. For a measurement of kaonic nitrogen transitions<sup>8</sup> nitrogen gas at 120 K 1.5 bar was filled instead of

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hydrogen. This measurement was done immediately before the hydrogen run and was also used to optimize the kaon degrader and as background measurement for the kaonic hydrogen data. Data corresponding to an integrated luminosity of  $32 \text{ pb}^{-1}$  ( $45.6 \cdot 10^6 \text{ K}^-$ ) were taken for nitrogen (272 GB files). In the final experiment we collected kaonic hydrogen data for  $58.4 \text{ pb}^{-1}$  ( $84.1 \cdot 10^6 \text{ K}^-$ , 288 GB files). Furthermore a background measurement with separated  $e^+ e^-$  beams and intentionally high X ray background was performed.

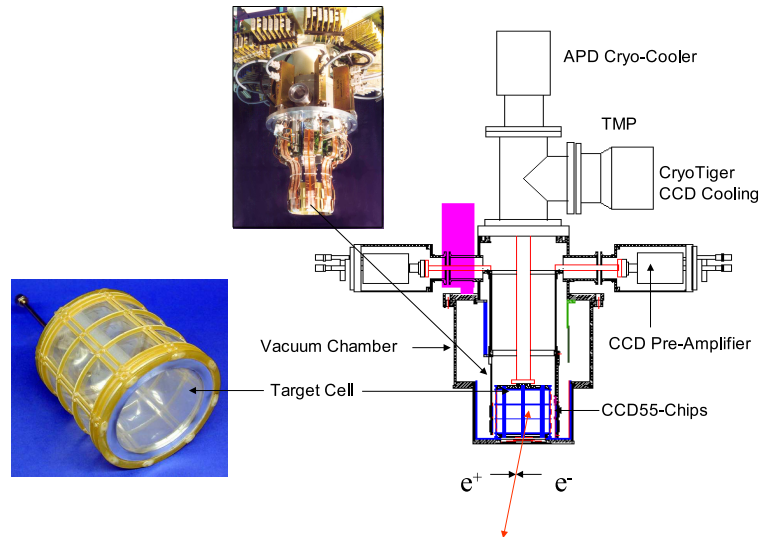


Fig. 2. Scheme of the experimental setup

The experimental challenge is the large low-energy X ray background (mainly from electron gamma showers resulting from lost  $e^+$  or  $e^-$  due to either Touschek scattering (interaction within the particle bunch) or interaction with residual gas. It was minimized by shielding of our experimental setup and by improvements in the beam optics achieved by the machine crew.

### 3. Data Analysis

The clean extraction of the kaonic X-ray spectrum requires careful treatment of the raw data. The procedure applied consists of the following steps:

The 'cluster analysis' selects single and double pixel events - the charge produced by minimal ionizing radiation is spread over a larger pixel area, producing a 'cluster' event which can be rejected. Even without irradiation CCDs have a charge

pedestal due to noise. This 'noise peak' is analyzed in position and width and data with broadened width are rejected. From position and width the noise threshold is derived: a pixel with charge above that threshold is regarded as a hit pixel. The CCD charge transfer efficiency depends on the position within the chip and is corrected. The individual exposures ('pictures') are added. For each detector the data are histogrammed and the individual energy calibrations using fluorescence background lines are calculated. The data of all individual detectors are added. A correction for the slightly different energy gains is applied. The overall resolution for the sum of detectors is determined.

The same procedure was applied to the data taken with the target filled with nitrogen ( $N$ -data), with hydrogen ( $H$ -data) and with hydrogen but separated beams which means without kaons ('No-collision' data). The data analysis for the nitrogen measurement is described in detail in <sup>8</sup>.

The spectra consist of a *continuous component* (low energy Bremsstrahlung and remaining events from minimal ionizing radiation) and *fluorescence lines* from excitation of structure- and detector materials and the *kaonic lines*.

The analysis of the kaonic hydrogen K lines meets 3 crucial challenges:

- The determination of the continuous background, since the width of the kaonic hydrogen lines is very sensitive to this.
- The iron  $K_\alpha$  X ray fluorescence line overlaps the kaonic  $K_\alpha$  signal. These lines must be disentangled.
- The yields of the higher kaonic hydrogen K transitions (5-1 and higher) can not be obtained directly from the data, the sensitivity of the result to theoretic input concerning the yield ratios has to be examined.

Two completely independent analysis starting from raw (pixel) data were done. In both procedures a  $H$  spectrum (see Fig.3) and a background spectrum with no kaonic hydrogen signal (denoted as  $B$ ) were fitted.

Approach 1 was to use for the background spectrum all 'No-collision' data and fit  $H$  and  $B$  data simultaneously by constructing a  $\chi^2$  function as the sum of 2 individual distributions  $\chi^2_H$  and  $\chi^2_B$  and use the same quadratic polynomial for the continuous background in  $H$  and  $B$ .

Four fit parameters for the continuous background were used: 1 for the normalisation of  $B$  vs.  $H$  plus 3 for the coefficients of a quadratic polynomial. The validity of using identical polynomial coefficients for  $H$  and  $B$  was checked by fitting the quotient spectrum  $H / B$  (excluding the KH signal regions) with a constant. For the selected fitrange of 2.9-10.4 keV no significant deviations were found. The electronic fluorescence lines from structure materials and the calibration foil (Ti) present in  $H$  and  $B$  were included in the fit model with the ratio given by the fitted normalization factor.

Approach 2 used the sum of  $N$  and a subset of the 'No-collision' data as background spectrum. This analysis proceeded in the following way: (1) Normalize the  $B$  spectrum to the same intensity in the Si fluorescence peak as the  $H$  spectrum. (2) Fit the normalized  $B$  spectrum using a cubic polynomial for the continuous

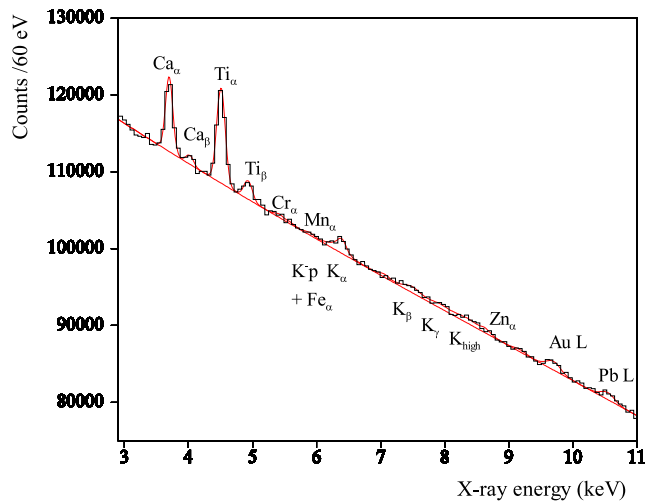


Fig. 3. X ray spectrum and fit curves. Nothing subtracted.

component and Gausspeaks for Cr, Mn, Ti, Fe, Zn, Au and Pb fluorescence lines and for the kaonic nitrogen lines at 4.58 keV and 7.59 keV. (3) Fit the  $H$  spectrum using a cubic polynomial for the continuous component and Gausspeaks for the fluorescence lines and the kaonic Carbon transition (6-5) and Voigt peaks for the kaonic hydrogen lines. The intensity of the Fe  $K_{\alpha}$  peak is constrained by the result of the fit of the normalized  $B$  spectrum from step 2.

The fit model for the kaonic hydrogen lines used Voigt functions with fixed Gaussian width for the detector resolution known from fluorescence lines and kaonic nitrogen transitions. Fit parameters are the intensities of  $K_{\alpha}$ ,  $K_{\beta}$  and  $K_{\gamma}$ , the energy of  $K_{\alpha}$  and the hadronic width of the kaonic hydrogen lines. One Lorentzian width was fitted and used for all K lines, since the hadronic broadening comes from the ground state width only. The energy *differences* of the kaonic hydrogen lines as compared to  $K_{\alpha}$  were fixed to the electromagnetic values.

Concerning the influence of the  $K_{high}$  transitions, we restricted the fitrange to a region where these contribute only by the effect of their low energy tails. For the individual  $K_{high}$  yields and  $K_{high}/K_{\gamma}$  yields we used values from theory<sup>9</sup> ( $K_{high}^{(1)}$ ,  $K_{high}^{(1)}/K_{\gamma}$ ). Furthermore we tested the effect of a different set of yields by assuming that the higher yields decrease exponentially and leaving this parameter float in a fit using the whole energy region. This procedure gave us an alternative set of yields,  $K_{high}^{(2)}$  and a value  $K_{high}^{(2)}/K_{\gamma}$ . Applying a restricted fitrange, as described above, we find that the results using the theoretical yield values and the

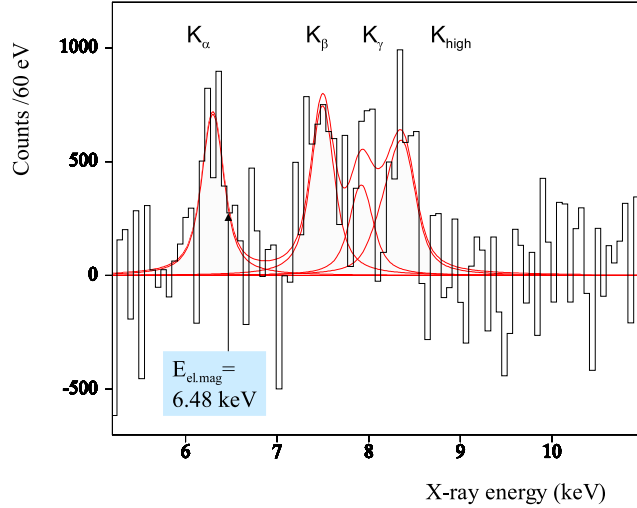


Fig. 4. Kaonic hydrogen X ray spectrum. All fit components except the kaonic lines are subtracted.

results using the alternative yields is consistent within  $\pm 3$  eV.

#### 4. Result and Conclusions

We obtained a transition energy of kaonic hydrogen  $K_\alpha$  of  $6287 \pm 37$  (stat.)  $\pm 6$  (syst.) eV. Using 6480 eV as transition energy without strong interaction, the  $1s$  strong interaction shift can be deduced:

$$\epsilon_{1s} = -193 \pm 37 \text{ (stat.)} \pm 6 \text{ (syst.) eV.} \quad (1)$$

We use here the so-called "Warsaw Convention" (measured X-ray minus calculated energy) for the sign of the shift. The negative sign of the shift means a repulsive strong interaction. The  $1s$  absorption (strong interaction) width (FWHM) result is

$$\Gamma_{1s} = 249 \pm 111 \text{ (stat.)} \pm 30 \text{ (syst.) eV.} \quad (2)$$

The systematic error covers detector energy calibration and energy resolution. In addition, systematic effects like fit range and method were also included.

The results are nearly compatible for the shift and compatible for the width of an earlier experiment<sup>3</sup> but our errors are significantly smaller (see Fig. 5). The values for the hadronic ground state shift and width are smaller than those from<sup>3</sup>. Our experimental results can test new theoretical studies<sup>10</sup>.

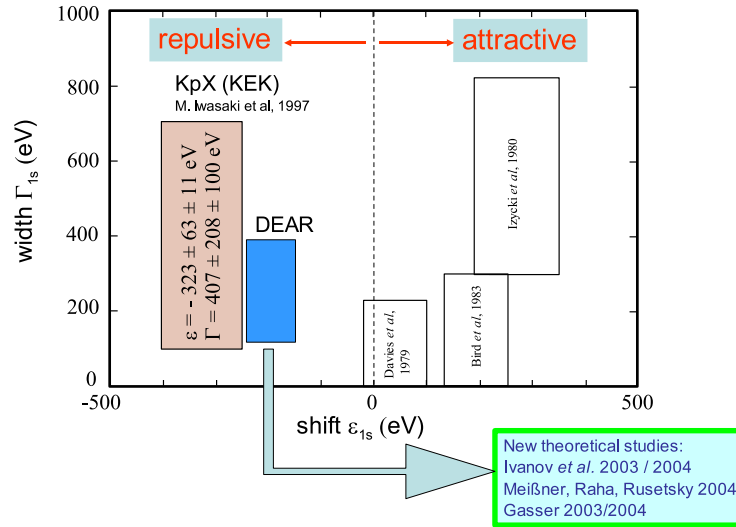


Fig. 5. The measured hadronic shift and width in kaonic hydrogen. Statistical and systematic errors have been added linearly.

### Acknowledgements

The DEAR collaboration wishes to thank the DAΦNE machine crew for their friendly collaboration. Part of the work was supported by the European Union "Transnational access to Research Infrastructure" (TARI) Contract No. HPRI-CT-1999-00088 and partially supported by the Swiss Science Foundation.

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