

A photon-in-cell code for laser wakefield accelerators

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Background

Nonlinear wave driven processes in plasmas are normally described by

➢A monochromatic pump wave coupling to other monochromatic waves

Random phase waves coupling to other random phase waves

Alternative:

➢ Random or broadband wave couples to monochromatic or coherent structures.

Introduction

► Wave kinetics: a novel approach to the interaction of waves with plasmas [1-5].

A monochromatic wave is described as a compact distribution of quasi-particles.

Broadband turbulence is described as a gas of quasiparticles (photons, plasmons, driftons, etc.)

[1] R. Trines et al., in preparation (2004).

- [2] J. T. Mendonça, R. Bingham, P. K. Shukla, Phys. Rev. E 68, 0164406 (2003).
- [3] L.O. Silva et al., IEEE Trans. Plas. Sci. 28, 1202 (2000).
- [4] R. Bingham et al., Phys. Rev. Lett. 78, 247 (1997).
- [5] R. Bingham et al., Physics Letters A 220, 107 (1996).

For complex phenomena

Photon acceleration/deceleration;
Photon wave breaking
Photon turbulence;
Modulational instabilities
Steepening/compression of short pulses

Applications

Applications in laser-plasma interaction...

Laser-plasma accelerators

Broadband laser-plasma interactions – Induced Spatial Incoherence

Strong plasma turbulence – Langmuir wave collapse

... and many other fields

Zonal flows in a "drifton gas" – anomalous transport driven by density/temperature gradients

Turbulence in planetary, stellar atmospheres – Rossby waves

> Plasma waves driven by neutrino bursts

➢The list goes on

Motivation

> Develop a PIC-type code for quasi-particle model of shortwavelength modes (photons, drift modes, ...) in a fluid plasma

Study of photon-driven wakefields for accelerators.

Basic features

Less computationally demanding than a full-PIC description;

A much simpler description than the "slowly evolving envelope approximation;"

Gives better physical insight into the EM field dynamics;

➢ Allows for an easy description of broadband, incoherent light pulses

Liouville theory

We define the quasiparticle density as the wave energy density divided by the quasiparticle energy:

$$N(\vec{k},\vec{r},t)=W(\vec{k},\vec{r},t)/\omega(k)$$

The number of quasiparticles is conserved:

$$\frac{d}{dt} \quad N(\vec{k}, \vec{r}, t) d\vec{r} d\vec{k} = 0$$

Then we can apply Liouville's theorem:

$$\frac{d}{dt}N(\vec{k}(t),\vec{r}(t),t) = -\frac{1}{t} + \vec{v} - \frac{1}{\vec{r}} + \vec{F} - \frac{1}{\vec{k}}N(\vec{k}(t),\vec{r}(t),t) = 0$$

where $\vec{v} = d\vec{x} / dt = \omega / \vec{k}$ and $\vec{F} = d\vec{k} / dt = -\omega / d\vec{x}$ are obtained from the q.p. dispersion relation

Condensed theory

Fluid model for the plasma (density n, momentum p):

$$\frac{n}{t} + \frac{np}{x} = 0 \qquad \frac{p}{t} + \frac{\gamma}{x} = -E \qquad \frac{E}{x} = 1 - n \qquad \gamma = \sqrt{1 + p^2 + A^2}$$

For a given photon density N(t,x,k), we have:

$$A^2 = \frac{N(t, x, k)}{\omega_k} dk$$

Particle model for the photons:

> Photon number conservation;

Hamiltonian: $\omega_i = \sqrt{\omega_p^2(t, x_i) + k_i^2}$

> Equations of motion:

$$\frac{dx_i}{dt} = \frac{k_i}{\omega_i} \qquad \qquad \frac{dk_i}{dt} = \frac{1}{2\omega_i} \frac{n}{x} \frac{n}{\gamma} (x_i, t)$$

Example: solitons (solitary pulses)

Soliton: a pulse exciting a relativistic plasma wave of only one bucket [6];

>All the pulse's photons are confined to that bucket;

Simultaneous acceleration and deceleration of photons keeps the pulse together;

Solitons are particularly well suited to be studied through the photon kinetic approach.

[6] P.K. Kaw, A. Sen, and T. Katsouleas, Phys. Rev. Lett. 68, 3172, 1992.

1-D Hybrid Code

- Monospaced grid, 3000 cells, up to 100 000 particles;
- Second order Runge-Kutta methods to solve both particle and fluid equations;
- Quadratic splines for photon and field projection;
- >A four-pass binomial filter on the vector potential;
- Moving simulation window: important for wakefield studies, Raman instabilities, etc.;
- Arbitrary shapes for the initial plasma and photon density profiles;
- The non-linear approach allows a peak plasma density of $n/n_0 = 10$ before breakdown occurs!

Results 1: long pulse

Modulational instability

A long, low-intensity laser pulse has been injected into an underdense plasma ($k_0/k_p = 7$);

Modulational instability clearly visible in all plots;

➢In addition, "wave breaking" of a photon wave, photon bunching, and photon turbulence, can be seen.



Results 2: single short pulse

Wakefield excitation

A short, low-intensity laser pulse has been injected into an underdense plasma ($k_0/k_p = 14$);

Plots of the pulse envelope show initial steepening...

>...followed by stretching and flattening of the pulse.



Single short pulse (cont'd)

➢ In photon phase space, there is initial photon deceleration, explaining the pulse steepening...

➤...while later on, the photon bunch is stretched and chirped, explaining the pulse stretching;

Clearly we are far removed from a soliton solution.



Results 3: double short pulse

► Non-linear soliton solutions often display a train of peaks in the pulse envelope...

> ...which motivated us to use a pulse with two very short peaks, crammed into a single plasma period;

➢ Pulse envelope plots already show that the pulse energy is confined better.



Double short pulse (cont'd)

➢ In photon phase space, we see that when one peak is decelerated, the other is accelerated...

…leading to a much more balanced energy exchange between photons and plasma...

≻and we obtain a mostly stable pulse:

Towards the soliton!



Conclusions

Photon kinetics...

► Is well suited for the study of interaction of EM waves with underdense plasmas,

➢ For 'old' phenomena: pulse steepening / compression and modulational instabilities,

➤And 'new' phenomena: photon acceleration, photon wave breaking and photon turbulence, soliton formation,...

Leads to improved physical insight through extended diagnostics,

Future work:

> 2-D/3-D simulations

> other types of fast waves (drift waves, Langmuir waves)

Photon kinetics: a valuable new perspective on wave-plasma interaction!